



Numerical Study of MHD Casson Liquid Stream over an Infinite Vertical Porous Plate with Newtonian Heating

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Abstract

Objectives: Magnetohydrodynamics (MHD) stream of Casson liquid over a limitless vertical porous plate with Newtonian heating is considered. Casson liquid is non-Newtonian liquid. **Methods/statistical analysis:** The governing nonlinear equations are comprehended utilizing shooting system. **Findings:** The impacts of relevant governing parameters on the liquid velocity and the temperature are exhibited graphically. The coefficient of skin friction and the rate of heat transfer are determined numerically. The present outcomes have been good simultaneousness with existing outcomes under some unique cases. **Applications:** Porous plates are additionally used in the design of heat exchangers, computer assemblies, polymer industry, and automotive industry.

Keywords: MHD, Casson Liquid, Newtonian Heating, Shooting Technique, Vertical Porous Plate.

Nomenclature

u' is the velocity of the fluid in x' -direction	β_1 is the volume coefficient of thermal expansion.
T' is the fluid temperature	B_0 is strength of the magnetic field along y' -axis.
h is the coefficient of heat transfer	β is the Casson parameter
k is the thermal conductivity	

<p>v' is the suction velocity ν is the kinematic viscosity T'_∞ is the free stream temperature g is the gravity of acceleration σ is the electrical conductivity of the fluid ρ is density of the fluid S is the suction parameter c_p is the specific heat constant q_r is the radiative heat flux along y'-axis</p>	<p>Gr is the thermal Grashof number M is the magnetic field parameter u and θ are the fluid flow velocity and the temperature respectively Pr is the Prandtl number R is the radiation parameter, Ec is the Eckert number. μ is the viscosity</p>
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1. Introduction

Magnetohydrodynamics (MHD) stream of Casson liquid over an infinite vertical porous plate is used in various Industrial and Engineering areas. Casson liquid is one of the non-Newtonian liquid. It is having profoundly viscosity nature. Human blood is one of the instances of Casson liquid. Porous plates are additionally used in the design of heat exchangers, computer assemblies, polymer industry, and automotive industry. The employments of magnetohydrodynamics are cooling of nuclear reactors, controlling boundary layer stream, geothermal energy extraction, sun power vitality gatherers, plasma aerodynamics, etc. The impact of Casson fluid flow over channels was addressed by the researchers [1–9]. In Refs. [10–11], researchers developed MHD natural convection flow past an impulsively started infinite vertical porous plate with Newtonian heating in the presence of radiation and effects of thermal radiation and mass diffusion on free convection flow near a vertical plate with Newtonian heating. The researchers [12–19] discovered the effect of MHD over a various geometries (vertical channel, vertical cylinder, stretching sheet).

The present assessment MHD stream of Casson liquid over an endless vertical permeable plate with Newtonian heating is examined. The governing non-linear differential equations are settled utilizing Ruge–Kutta fourth-order strategy along with shooting system. The impacts of governing parameters such as Prandtl number, Magnetic field parameter, Casson parameter, Suction parameter, and radiation parameter on the liquid velocity and the temperature are exhibited in graphically. The skin friction coefficient and the Nusselt number are found in numerically and showed up in tabular structure.

2. Mathematical Formulation of the Problem

Consider the steady hydrodynamic stream of Casson liquid over a limitless vertical plate with Newtonian heating. The plate in the upward way along the x -axis and y -axis normal to it as showed up in Figure 1. Assumed that the heat transfer rate from the plate with a finite heat capacity is corresponding to the local surface temperature (T').

The under Boussinesq estimate for the continuity equation, conservation of mass, and the energy equations are as per the following:

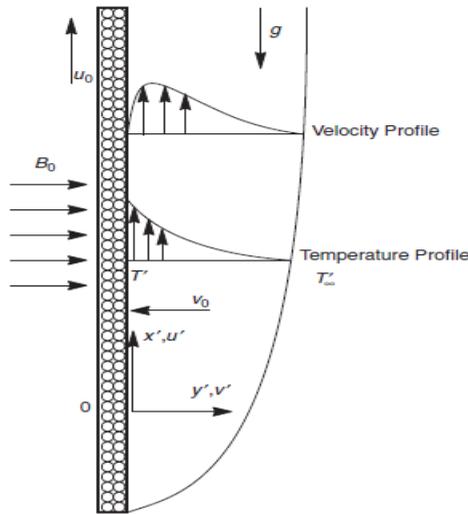


FIGURE 1. Physical model of the problem.

$$\frac{\partial v'}{\partial y'} = 0 \quad (1)$$

$$v' \frac{\partial u'}{\partial y'} = v' \left(1 + \frac{1}{\beta} \right) \frac{\partial^2 u'}{\partial y'^2} + g\beta_1 (T' - T_\infty') - \frac{\sigma B_0^2}{\rho} u' \quad (2)$$

$$v' \frac{\partial T'}{\partial y'} = \frac{k}{\rho c_p} \frac{\partial^2 T'}{\partial y'^2} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y'} + \frac{\mu}{\rho c_p} \left(1 + \frac{1}{\beta} \right) \left(\frac{\partial u'}{\partial y'} \right)^2 + \frac{\sigma B_0^2}{\rho} u'^2 \quad (3)$$

The limit conditions are as follows:

$$\left. \begin{aligned} u' &= u_0, \frac{\partial T'}{\partial y'} = -\frac{h}{k} T' \text{ at } y' = 0 \\ u' &\rightarrow 0, T' \rightarrow T_\infty' \text{ as } y' \rightarrow \infty \end{aligned} \right\} \quad (4)$$

Integrating equation (1) for suction constant, we have

$$v' = -v_0 \quad (5)$$

Where $v_0 > 0$ is the normal velocity of suction at the plate?

Now equations (2) and (3) are as the following:

$$-v_0 \frac{\partial u'}{\partial y'} = v' \left(1 + \frac{1}{\beta} \right) \frac{\partial^2 u'}{\partial y'^2} + g\beta_1 (T' - T_\infty') - \frac{\sigma B_0^2}{\rho} u' \quad (6)$$

$$-v_0 \frac{\partial T'}{\partial y'} = \frac{k}{\rho c_p} \frac{\partial^2 T'}{\partial y'^2} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y'} + \frac{\mu}{\rho c_p} \left(1 + \frac{1}{\beta} \right) \left(\frac{\partial u'}{\partial y'} \right)^2 + \frac{\sigma B_0^2}{\rho} u'^2 \quad (7)$$

The radiation heat flux q_r in equation (7) fulfills the accompanying non-linear differential equation [7]

$$\frac{\partial q_r}{\partial y'} = 4\alpha\sigma^* (T'^4 - T_\infty'^4) \tag{8}$$

where α is the absorption coefficient and σ^* is the Stefan–Boltzman constant. Assuming that the temperature contrasts inside the stream are adequately little such that T'^4 can be expressed as a linear function of T' and utilizing the Taylor series expression about $T_\infty'^4$ in the wake of disregarding higher order terms, we have:

$$T'^4 \cong 4T_\infty'^3 T' - 3T_\infty'^4 \tag{9}$$

From equations (8), (9), and (7) we get

$$-v_0 \frac{\partial T'}{\partial y'} = \frac{k}{\rho c_p} \frac{\partial^2 T'}{\partial y'^2} - \frac{16\alpha\sigma^* T_\infty'^3}{\rho c_p} (T' - T_\infty') + \frac{\mu}{\rho c_p} \left(1 + \frac{1}{\beta}\right) \left(\frac{\partial u'}{\partial y'}\right)^2 + \frac{\sigma B_0^2}{\rho} u'^2 \tag{10}$$

The non-dimensional quantities are as per the following:

$$\left. \begin{aligned} y = \frac{y'h}{k}, u = \frac{u'}{u_0}, \theta = \frac{T' - T_\infty'}{T_\infty'}, S = \frac{v_0 k}{h\nu}, M = \frac{\sigma B_0^2 k^2}{\mu h^2} \\ Gr = \frac{g\beta_1 T_\infty' k^2}{\nu u_0 h^2}, Pr = \frac{\mu c_p}{k}, Ec = \frac{\mu u_0^2}{k T_\infty'}, R = \frac{16\alpha\sigma^* k T_\infty'^3}{h^2} \end{aligned} \right\} \tag{11}$$

Using equation (11), from equations (6), (10), and (4) are as per the following:

$$\left(1 + \frac{1}{\beta}\right) \frac{d^2 u}{dy^2} + S \frac{du}{dy} + Gr\theta - Mu = 0 \tag{12}$$

$$\frac{d^2 \theta}{dy^2} + Pr S \frac{d\theta}{dy} - R\theta + Ec \left(1 + \frac{1}{\beta}\right) \left(\frac{du}{dy}\right)^2 + EcMu^2 = 0 \tag{13}$$

The limit conditions are

$$\left. \begin{aligned} u = 1, \frac{d\theta}{dy} = -(\theta + 1) \text{ at } y = 0 \\ u \rightarrow 0, \theta \rightarrow 0 \text{ as } y \rightarrow \infty \end{aligned} \right\} \tag{14}$$

Physical amounts of intrigue are the skin friction coefficient C_f and the local rate of heat transfer coefficient Nu are characterized as

$$C_f = \frac{\tau'}{\rho G r u_0^2} = -\left(1 + \frac{1}{\beta}\right) \frac{du}{dy} \Big|_{y=0} \text{ and } Nu = -\frac{\nu}{u_0 (T' - T_\infty')} \frac{dT'}{dy} \Big|_{y=0} = 1 + \frac{1}{\theta(0)} \tag{15}$$

3. Results and Discussion

The present assessment reveals that hydromagnetic stream of Casson liquid over a vast vertical porous plate with Newtonian heating. Here we explored the present outcomes for both Newtonian ($\beta \rightarrow \infty$) and non-Newtonian ($\beta = 0.5$) cases. Here the Casson parameter taken as $\beta = 0.5$. The nonlinear governing equations can be unwound utilizing Runge–Kutta fourth-order methodology along with shooting system. The impacts of governing parameters on the liquid velocity and the temperature distribution are appeared in graphically while the skin friction coefficient and the Nusselt number are determined numerically and are showed up in tabular structure. The present outcomes have been incredible simultaneousness with the current examinations under some restricted cases.

The impact of magnetic field parameter M on the liquid flow velocity $u(y)$ is showed up in Figure 2. We saw that the liquid velocity $u(y)$ lessens for higher estimations of magnetic field parameter M further the momentum boundary layer thickness diminishes this causes the Lorentz force. The effect of Suction parameter S on the liquid velocity $u(y)$ and the temperature distribution $\theta(y)$ are appeared in Figures 3 and 4. We saw that the liquid velocity $u(y)$ and the temperature $\theta(y)$ rot with the impact of S . This reality that of the impact of the suction is to take away the hot liquid on the vertical plate and consequently diminishing the liquid velocity with lessens free convection currents intensity. The variation on fluid velocity $u(y)$ with the impact of thermal Grashof number Gr is shown in Figure 5. We have seen that the liquid velocity $u(y)$ improves with growing

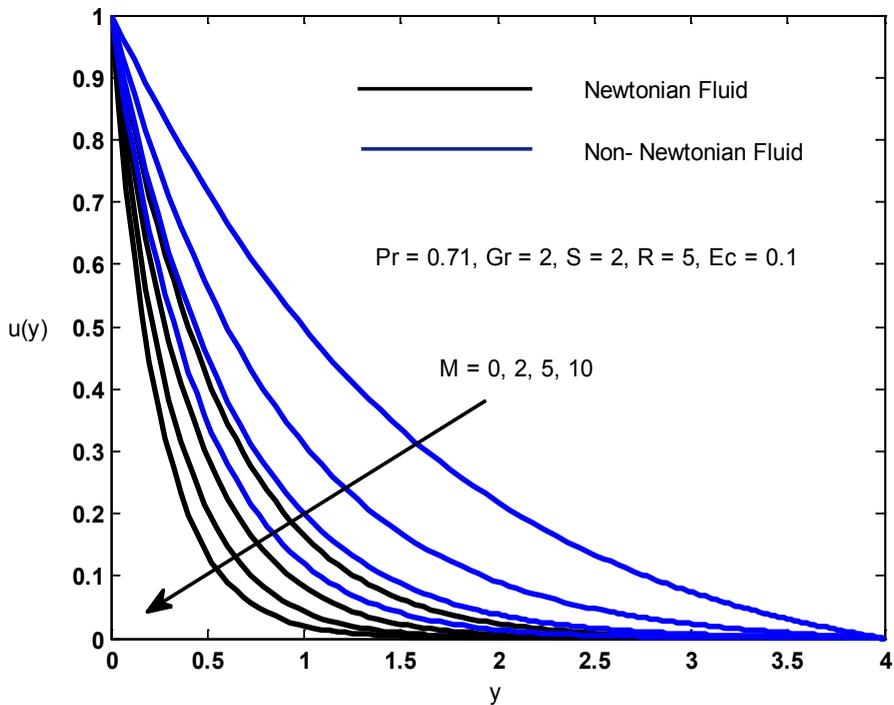


FIGURE 2. The fluid velocity $u(y)$ with M .

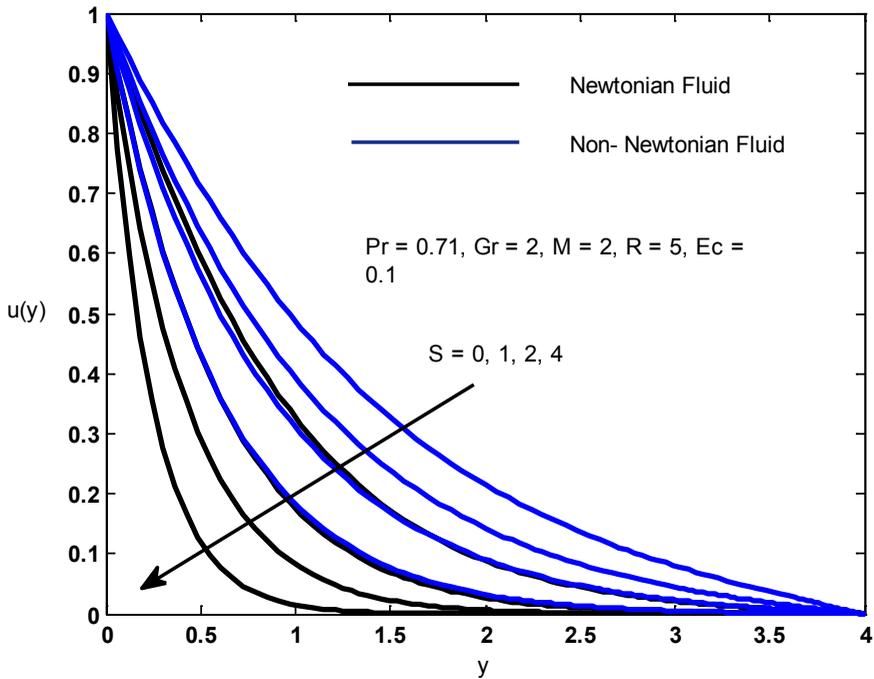


FIGURE 3. The fluid velocity $u(y)$ with S .

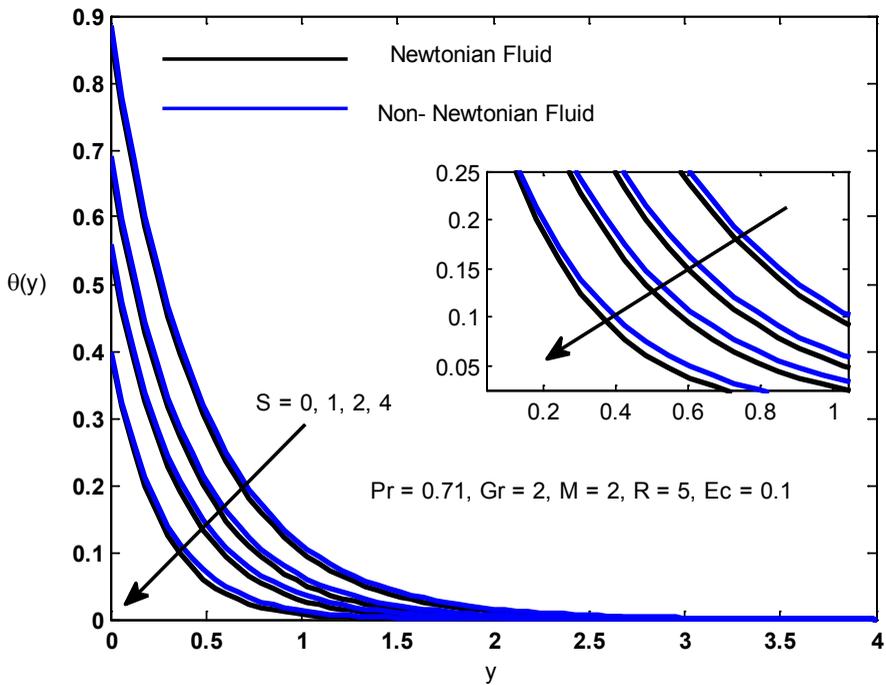


FIGURE 4. The temperature distribution $\theta(y)$ with S .

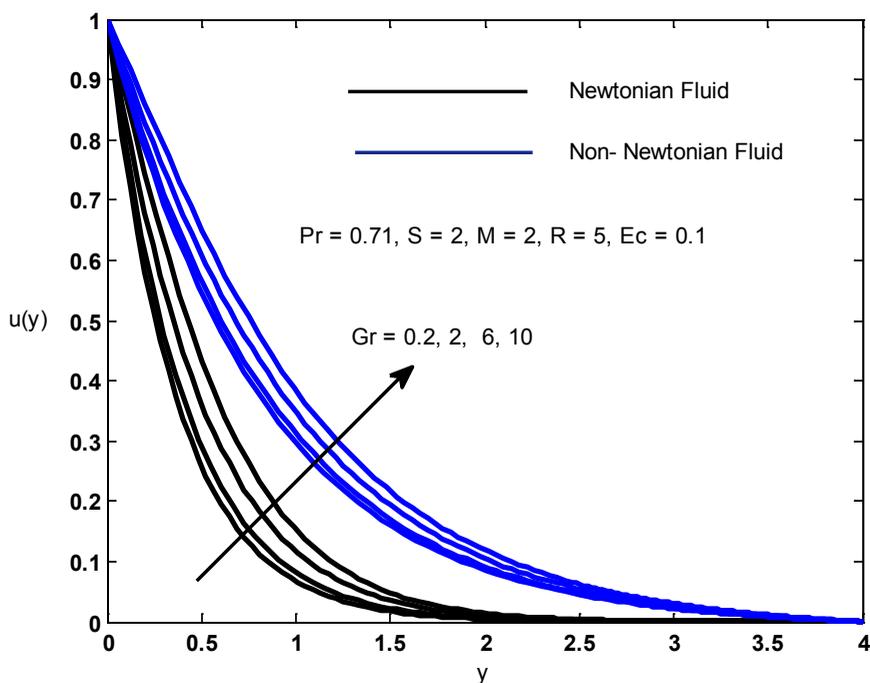


FIGURE 5. The fluid velocity $u(y)$ with Gr .

Grashof number. This causes the higher estimations of Grashof number then the viscous force impact rotates and additions in momentum boundary layer thickness. The impact of Casson parameter β on the liquid velocity $u(y)$ is demonstrated in Figure 6. We saw that the liquid velocity diminishes for higher estimations of Casson parameter. The yield pressure rotates which causes the creation for opposition force which make the liquid velocity decreases. The liquid is closer to the Newtonian liquid.

Figure 7 speaks to the impact of Prandtl number Pr on the temperature distribution $\theta(y)$. We uncover that the temperature rotates with extending the Prandtl number. This causes that for higher estimations of Prandtl number relate to decreases the thermal diffusivity. Figure 8 portrays the impact of radiation parameter R on the temperature distribution $\theta(y)$. We have seen that the temperature diminishes with a growing radiation parameter which causes for higher estimations of thermal radiation parameter have affinity to absorb the radiative heat and subsequently the buoyancy force rotates as well as thermal boundary layer lessens.

The numerical estimations of the skin friction coefficient and the Nusselt number for distinct governing parameters are showed up in Table 1. We saw that the skin friction coefficient overhauls for higher estimations of Prandtl number, magnetic field parameter, suction parameter and radiation parameter and the skin friction coefficient rotates with growing Grashof number and Eckert number. The Nusselt number improves for upgrading Prandtl number, Grashof number, radiation parameter and suction parameter and the Nusselt number declines for higher estimations of Eckert number and magnetic

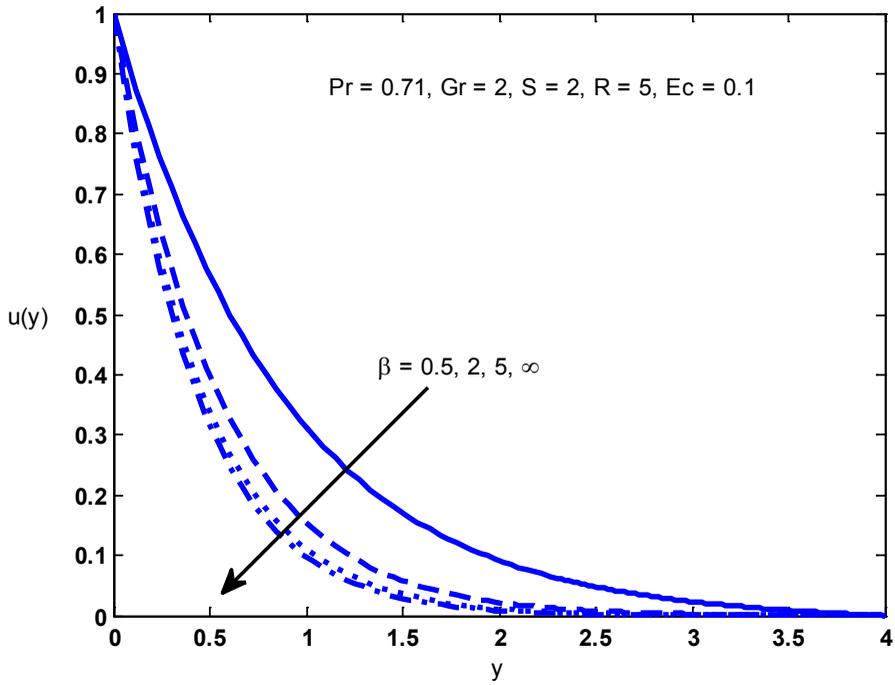


FIGURE 6. The fluid velocity $u(y)$ with β .

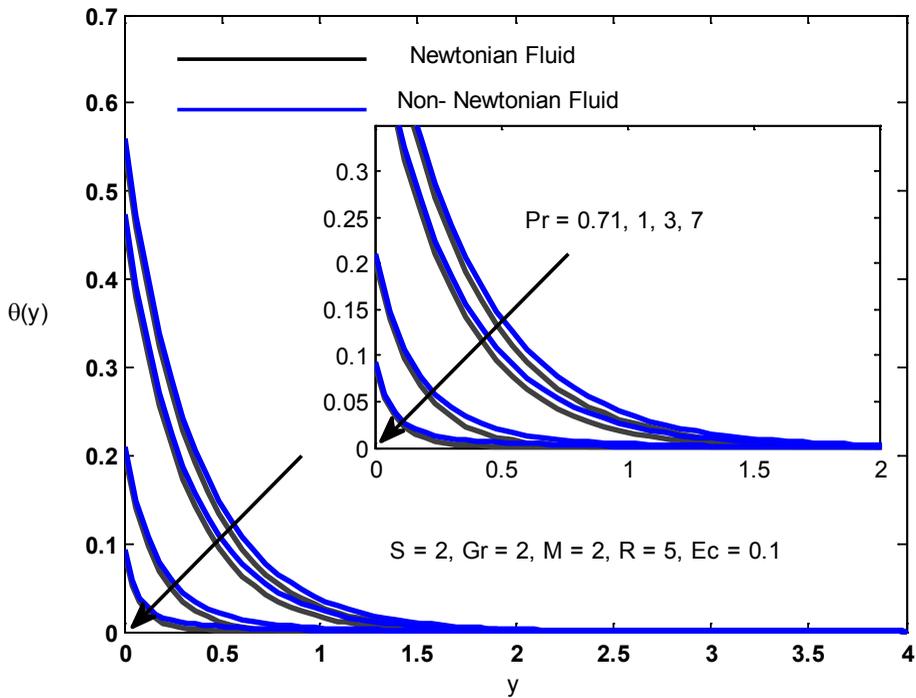


FIGURE 7. The temperature distribution $\theta(y)$ with Pr .

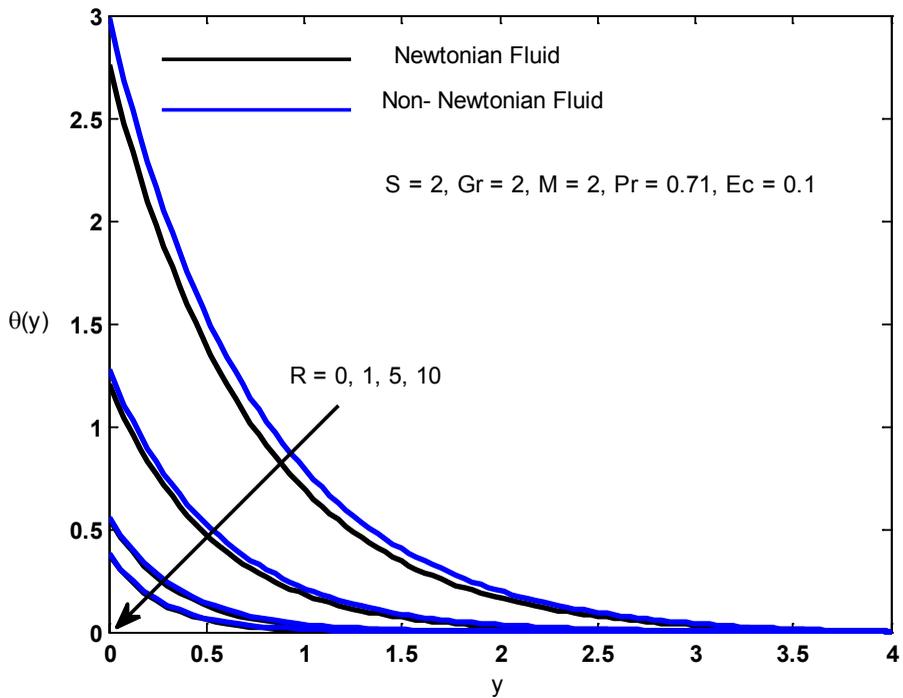


FIGURE 8. The temperature distribution $\theta(y)$ with R .

TABLE 1. The numerical values of the skin friction coefficient Cf and the rate of heat transfer Nu for distinct values of Pr, Gr, Ec, S, M , and R when the absence of $(\beta \rightarrow \infty)$

Pr	Gr	Ec	S	M	R	Cf		Nu	
						In [10]	Present result $(\beta \rightarrow \infty)$	In [10]	Present result $(\beta \rightarrow \infty)$
0.3	2	0.1	2	2	5	2.18782	2.18782	2.40855	2.40855
0.71						2.31881	2.31881	2.8436	2.8436
1						2.38021	2.38021	3.18447	3.18447
3						2.43219	2.43219	6.04361	6.04361
7						2.44549	2.44549	12.7096	12.7096
0.71	0.2					2.61649	2.61649	2.82578	2.82578
	6					1.80347	1.80347	2.86364	2.86364
	10					1.24627	1.24627	2.88043	2.88043
	15					0.53924	0.53924	2.88994	2.88994
	2	0				2.39937	2.39937	3.05598	3.05598
		0.05				2.38081	2.38081	2.94103	2.94103
		0.1				2.36233	2.36233	2.8392	2.8392

		0.15				2.34392	2.34392	2.74838	2.74838
		0.2				2.3256	2.3256	2.66687	2.66687
		0.1	0			0.93886	0.93886	2.15346	2.15346
			1			1.58837	1.58837	2.48615	2.48615
			4			3.91856	3.91856	3.59115	3.59115
			6			5.60228	5.60228	4.33171	4.33171
			2	0		1.63506	1.63506	2.94084	2.94084
				5		2.99838	2.99838	2.74832	2.74832
				10		3.43956	3.43956	2.66563	2.66563
			2	0		0.99746	0.99746	1.4348	1.4348
				1		1.7479	1.7479	1.83538	1.83538
				5		2.05239	2.05239	2.85988	2.85988
				10		2.14835	2.14835	3.69343	3.69343
				20		2.21501	2.21501	4.90493	4.90493

field parameter. The present outcomes have been great simultaneousness with existing outcomes [6] when the nonattendance of Casson parameter ($\beta \rightarrow \infty$).

4. Conclusions

In this study, we uncover that the hydromagnetic stream of Casson liquid over an infinite vertical porous channel with Newtonian heating is examined. The nonlinear differential equations are grasped utilizing shooting strategy. The present outcomes are contrasted differently in relation to Newtonian and non-Newtonian liquid and are demonstrated graphically. The impact of the suitable governing parameters on the liquid velocity and the temperature are showed up in graphically. The skin friction coefficient and the rate of heat transfer are surveyed numerically and showed up in tabular structure. The destinations of the present outcomes are as follows:

- The momentum boundary layer diminishes with the impact of magnetic parameter. The liquid velocity and the temperature are decay for higher estimations of suction parameter.
- The momentum boundary layer increment for higher estimations of Grashof number and inverse conduct is seen with extending Casson parameter.
- Thermal boundary layer rots for higher estimations of Prandtl number and Radiation parameter. The skin friction coefficient increments with extending the Prandtl number, magnetic field parameter, suction parameter and radiation parameter and the skin friction coefficient diminishes for higher estimations of Grashof number and Eckert number. The Nusselt number improves for upgrading Prandtl number, Grash of

number, radiation parameter and suction parameter and the Nusselt number decreases for higher estimations of Eckert number and magnetic field parameter.

- The present outcomes have been incredible simultaneousness with existing outcomes [10] when the nonappearance of Casson parameter i.e. ($\beta \rightarrow \infty$).

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